Geodesic motion and kludge waveforms from a test-mass orbiting around 'quasi-Kerr' object.

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- provide evidence in favour of GR's 'no-hair' theorem (uniqueness of Kerr metric)
- reveal the true identity of the 'dark objects' in galactic nuclei (Kerr BH vs Boson stars, ...)

Building 'quasi-Kerr' metric

Start with the exterior Hartle-Thorne metric describing the space-time of any axisymmetric & stationary body up to $\mathcal{O}(J^2)$ accuracy:

$$g_{\alpha\beta}^{HT} = g_{\alpha\beta}^{HT,Kerr}(up \, to \, \mathcal{O}(J^2)) + \epsilon h^{HT} + \mathcal{O}(J^{3+}),$$

where
$$\epsilon = -Q/M^3 - (J/M^2)^2$$
. For Kerr BH $\epsilon = 0$,
 $Q/M^3 = -(J/M^2)^2$, for NS $Q/M^3 = -\alpha (J/M^2)^2$, where $\alpha \in [2, 12]$.

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Transform to Boyer-Lindquist coordinates and

$$g_{\alpha\beta}^{QK} = g_{\alpha\beta}^{Kerr} + \epsilon h_{\alpha\beta}$$

Depends on M, $a = J/M^2$, ϵ .

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- Non-separable for generic orbits, lack of Carter constant.
- Employ canonical perturbation theory: H = H₀ + ϵH₁.
 Constants of motion {Q₀, P₀} at zero-order, now should satisfy (at first order)

$$\dot{Q} = -\epsilon \left(\frac{\partial H_1}{\partial P}\right)_0 \quad \dot{P} = \epsilon \left(\frac{\partial H_1}{\partial Q}\right)_0$$

Isolate secular changes by time averaging.

$$P = \langle \dot{P} \rangle t + P_0, \quad Q = \langle \dot{Q} \rangle t + Q_0$$

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Study periastron shift for Kerr and 'quasi-Kerr' orbits: $\Delta \phi_{Kerr} - \Delta \phi_{qK}$, where

$$\Delta \phi = \int_0^{2\pi} \phi(\chi) \ d\chi.$$

Periastron shift 1



Periastron shift 2



Periastron shift 3



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• Compute overlaps between Kerr and 'quasi-Kerr' waveforms without radiation reaction, but truncated at T_{RR} defined above.

Comparison of waveforms



Overlaps





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