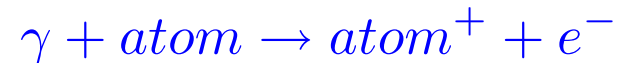


Photoelectric absorption

Photoelectric absorption

A bound electron can absorb completely the energy of a photon :



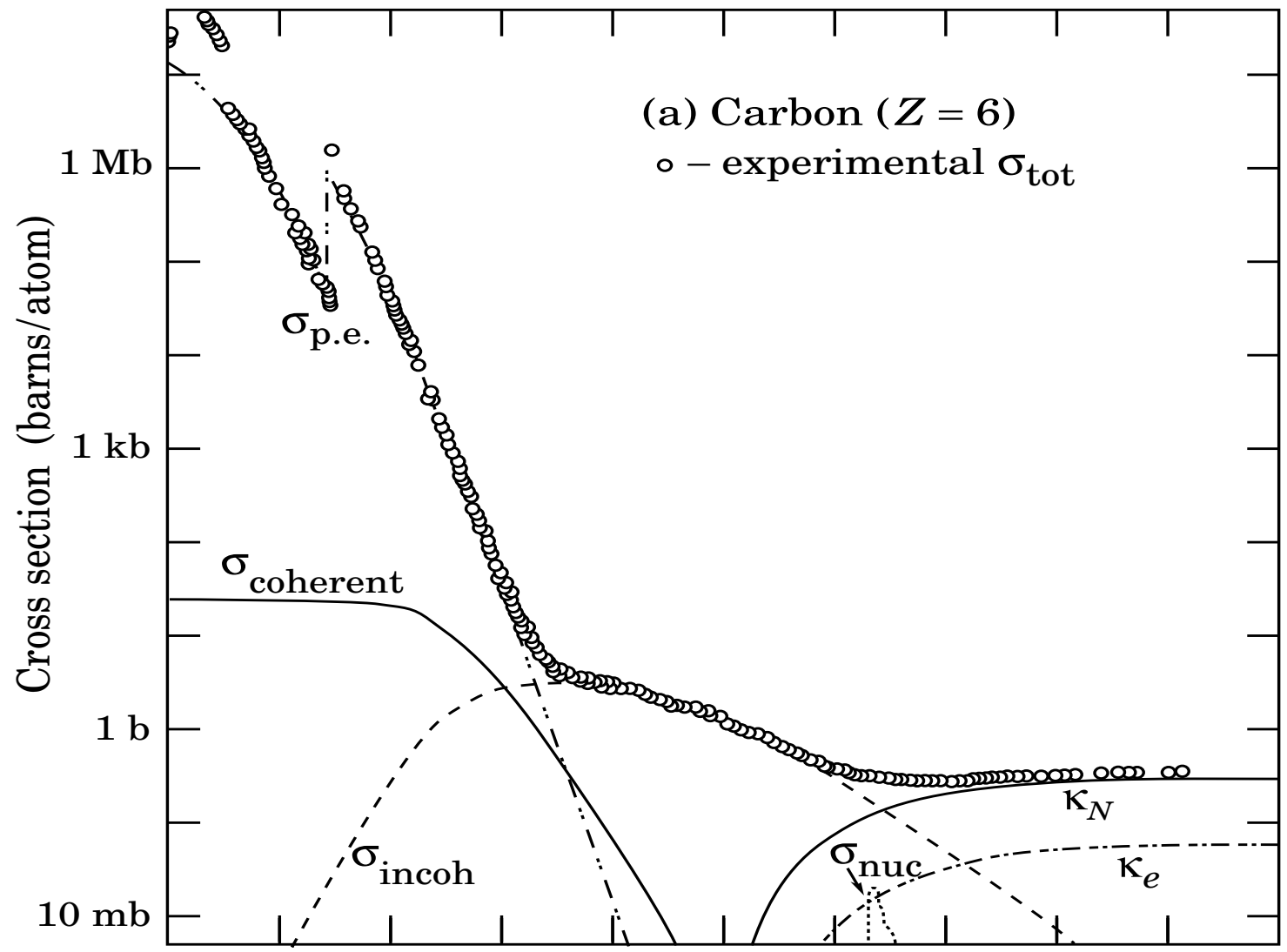
The electron is ejected with kinetic energy $T = E_\gamma - B_s$.

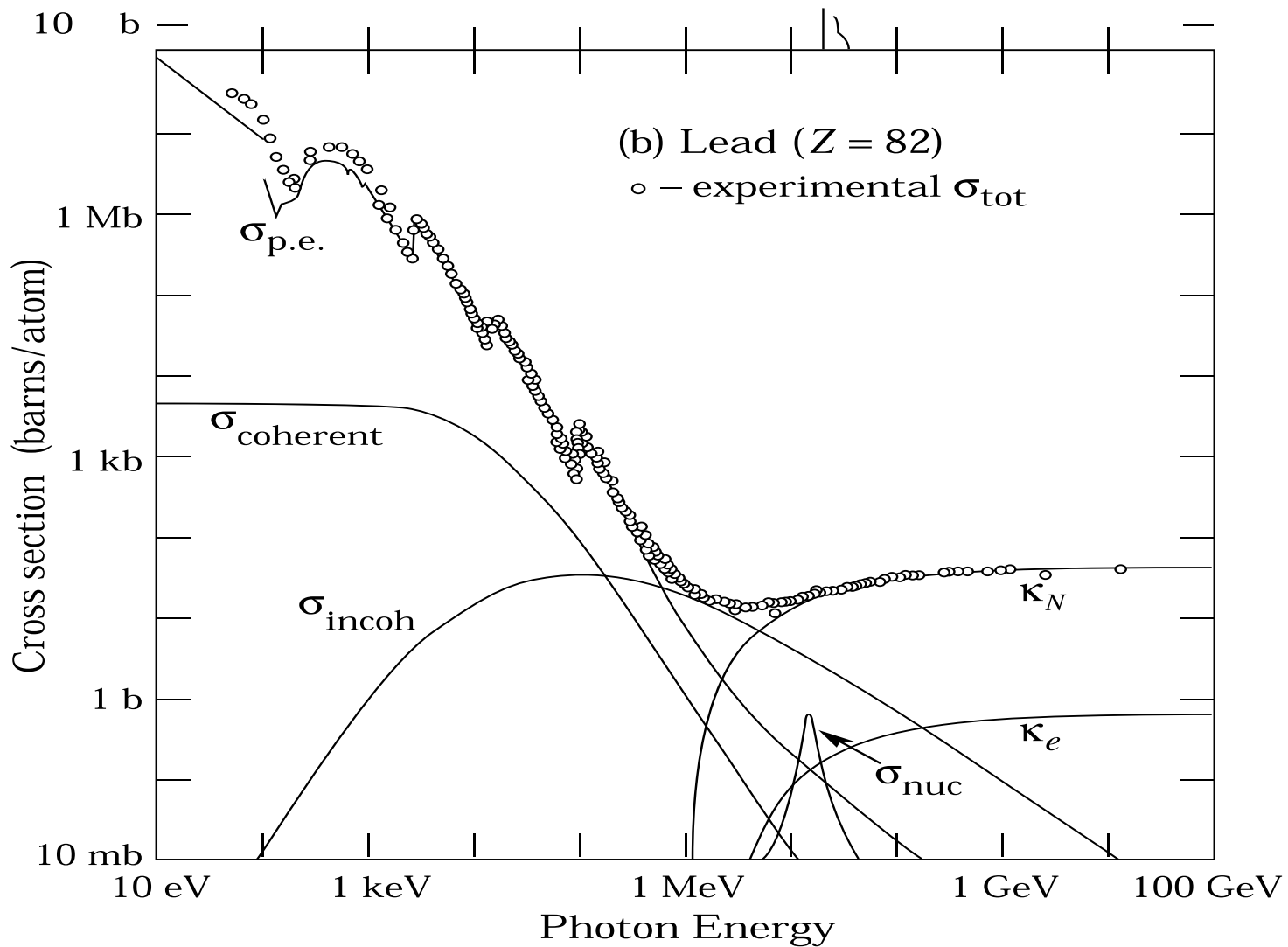
(E_γ : energy of the incident photon, B_s : binding energy of the corresponding subshell). The nucleus absorbs the recoil momentum.

The cross section per shell can be parametrized [Biggs87] :

$$\sigma_s = r_e^2 \alpha^4 Z^5 f \left(\frac{1}{E_\gamma^{a(E_\gamma)}} \right)$$

with f : nonsimple function of $1/E_\gamma$, and $1 \leq a(E_\gamma) \leq 4$.





The total cross section has discontinuities at $E_\gamma = B_s$
(absorption edge).

There are several parametrizations and tables of the cross sections.
See [Cullen97, Biggs87].

If $E_\gamma > B_K$ the absorption occurs mainly on the K-shell (80% of the cases).

The electron is emitted forward in the direction of the incident photon at high E_γ , and perpendicular to the photon at low E_γ [Sauter31, Gavri61].

Following the photoabsorption in the K-shell, characteristic X-rays or Auger electrons are emitted [Perkin91].

Cross section per atom

We use a parametrisation of the photoabsorption cross section proposed by Biggs and al. [Biggs87] :

$$\sigma(Z, E_\gamma) = \frac{a(Z, E_\gamma)}{E_\gamma} + \frac{b(Z, E_\gamma)}{E_\gamma^2} + \frac{c(Z, E_\gamma)}{E_\gamma^3} + \frac{d(Z, E_\gamma)}{E_\gamma^4} \quad (1)$$

The coefficients a, b, c, d are fitted with experimental data by the least square method separately in each energy interval [Grichi94].

As a rule, the interval borders are equal to the corresponding photoabsorption edges.

Mean free path

$$\lambda(E_\gamma) = \left(\sum_i n_{ati} \cdot \sigma(Z_i, E_\gamma) \right)^{-1}$$

n_{ati} : nb of atoms per volume of the i^{th} element in the material.

At initialization stage, the function `BuildPhysicsTables()` computes and tabulates :

- `crossSectionPerAtom` for all elements
- `meanFreePath` for all materials

The cross section and mean free path can be tabulated only above 50 keV. Below this, they are too discontinue: they are recomputed 'on fly' from formula 1.

final state

choose an Element : the binding energy of the shells depend of Z_i . In a compound material one choose randomly an Element according :

$$Prob(Z_i, E_\gamma) = \frac{n_{ati}\sigma(Z_i, E_\gamma)}{\sum_i [n_{ati} \cdot \sigma_i(E_\gamma)]}$$

final state : the simulation is presently rather crude.

A quanta can be absorbed if $E_\gamma > B_{shell}$. The shell energies are taken from `G4AtomicShells` data. One choose the **closest** atomic shell available.

The photoelectron is emitted with kinetic energy:

$$T_{photoelectron} = E_\gamma - B_{shell}(Z_i)$$

Theta distribution of the photoelectron

The polar angle of the photoelectron is sampled from the Sauter-Gavrila distribution (for K-shell) [Gavri59], which is correct only to zero order in αZ :

$$\frac{d\sigma}{d(\cos \theta)} \sim \frac{\sin^2 \theta}{(1 - \beta \cos \theta)^4} \left\{ 1 + \frac{1}{2} \gamma(\gamma - 1)(\gamma - 2)(1 - \beta \cos \theta) \right\} \quad (2)$$

where β and γ are the Lorentz factors of the photoelectron.

Relaxation

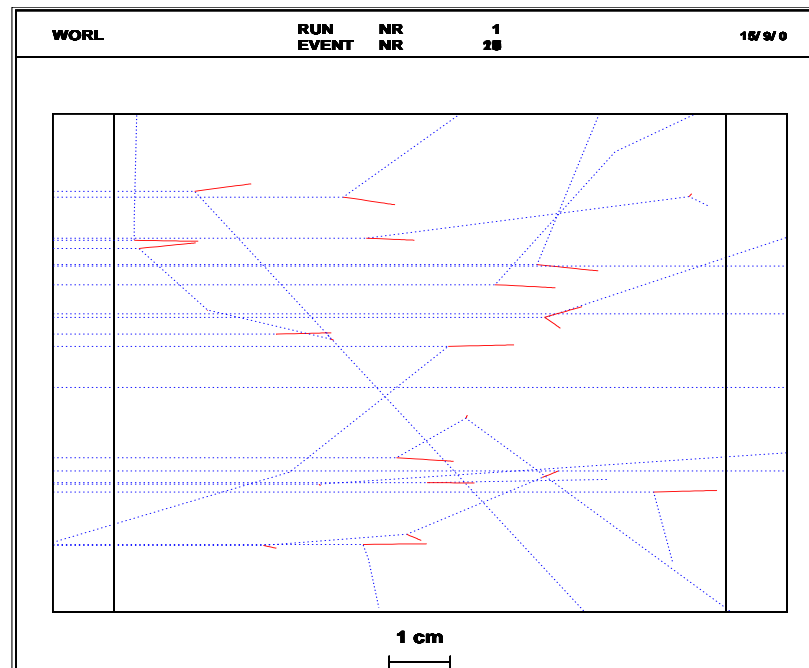
In the current implementation the relaxation of the atom is not simulated, but instead is counted as a local energy deposit.

attenuation

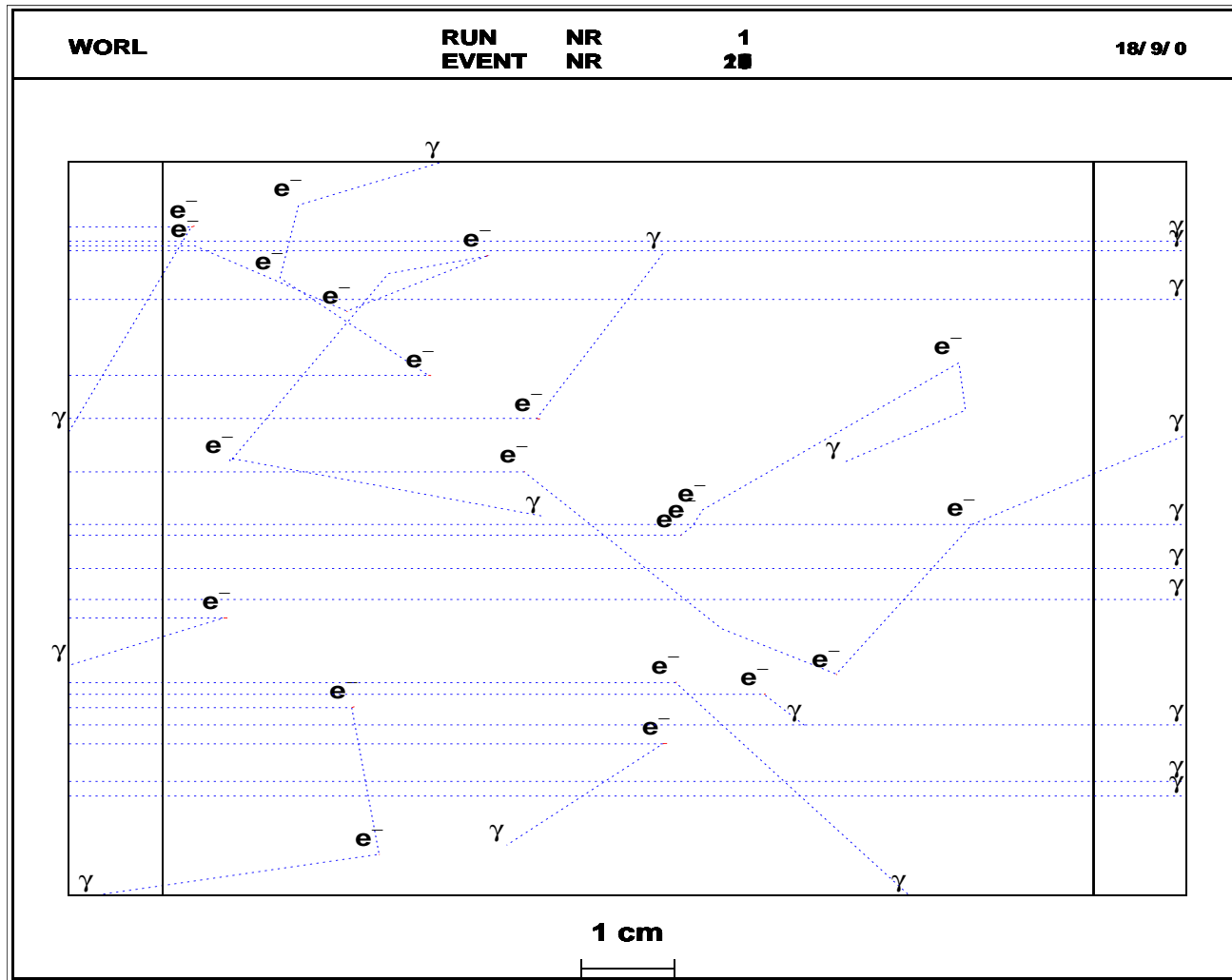
$$\sigma_{tot} = \sigma_{pair} + \sigma_{comp} + \sigma_{phot} + \sigma_{rayl} \quad \longrightarrow \quad \mu = n_{at} \sigma_{tot}$$

A beam of monoenergetic photons is attenuated in intensity (not in energy) according to : $I(x) = I(0) \exp(-\mu x) = I(0) \exp(-x/\lambda)$

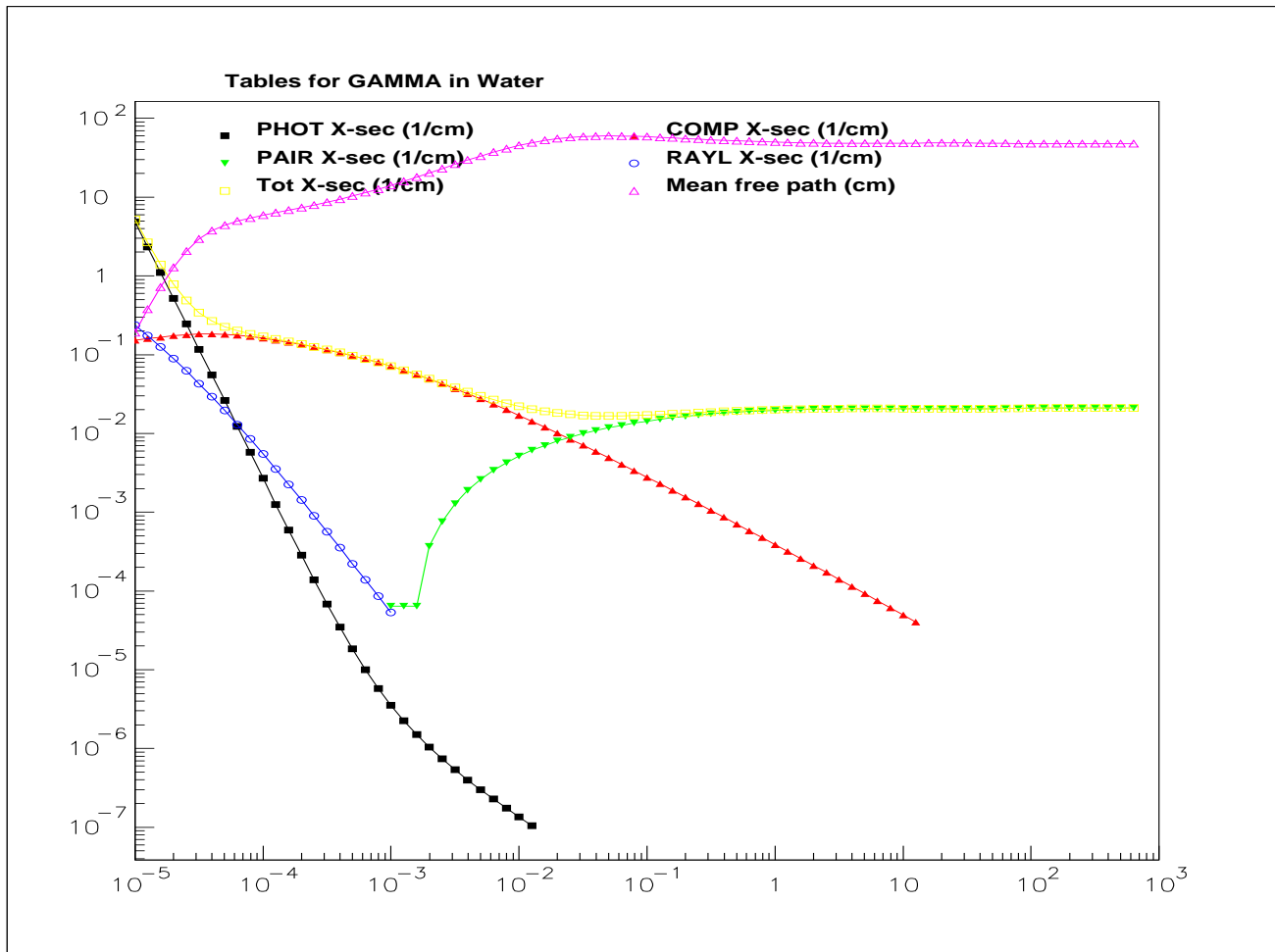
Below : 20 photons, 5 MeV, entering 10 cm of Al. 4 exit unaltered.



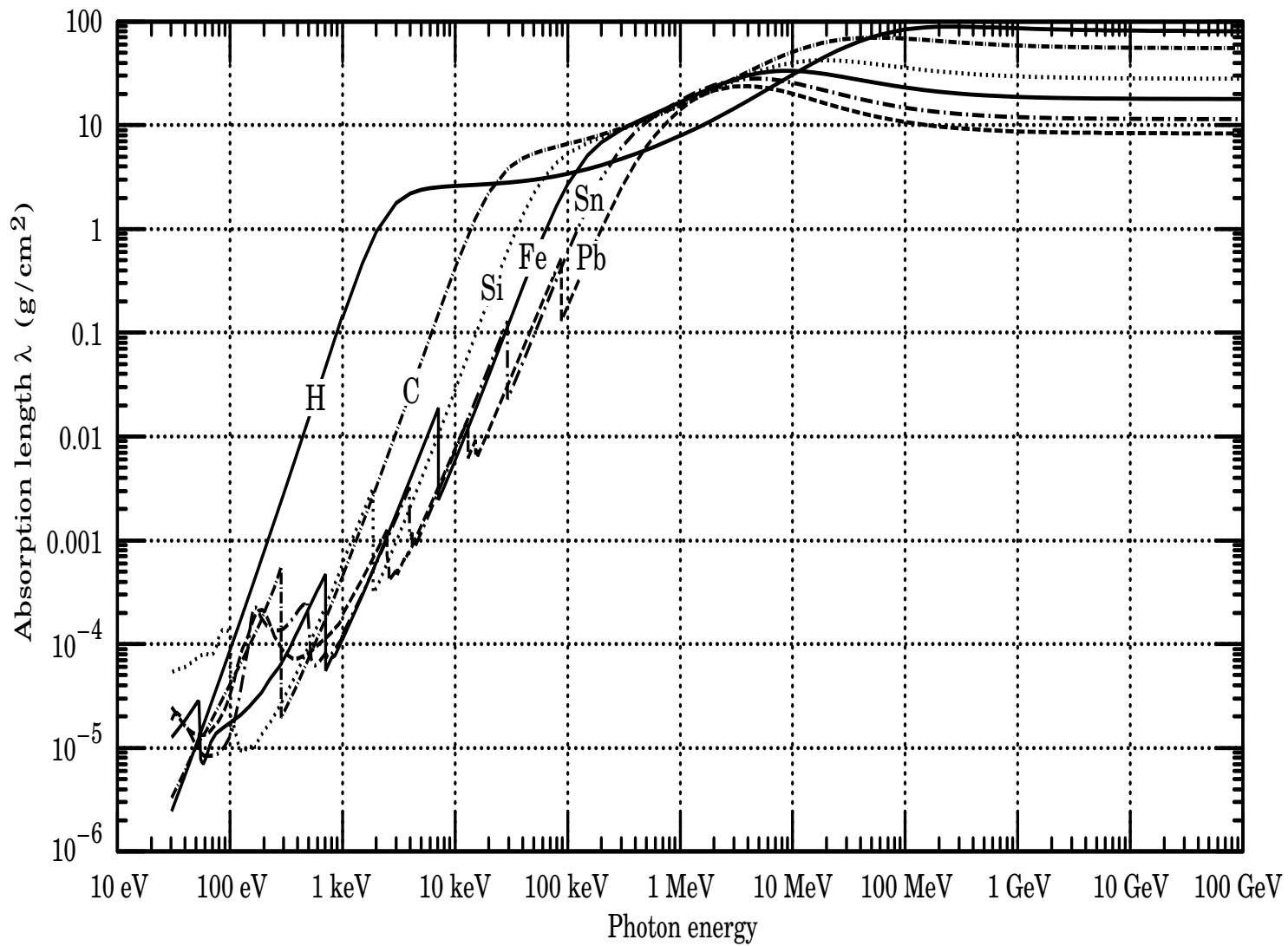
20 photons, 400 keV, entering 10 cm of water.
 (compare with e^- and protons)



Macroscopic cross sections for photon in water. (→ mean free path)



photon energy (GeV)



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