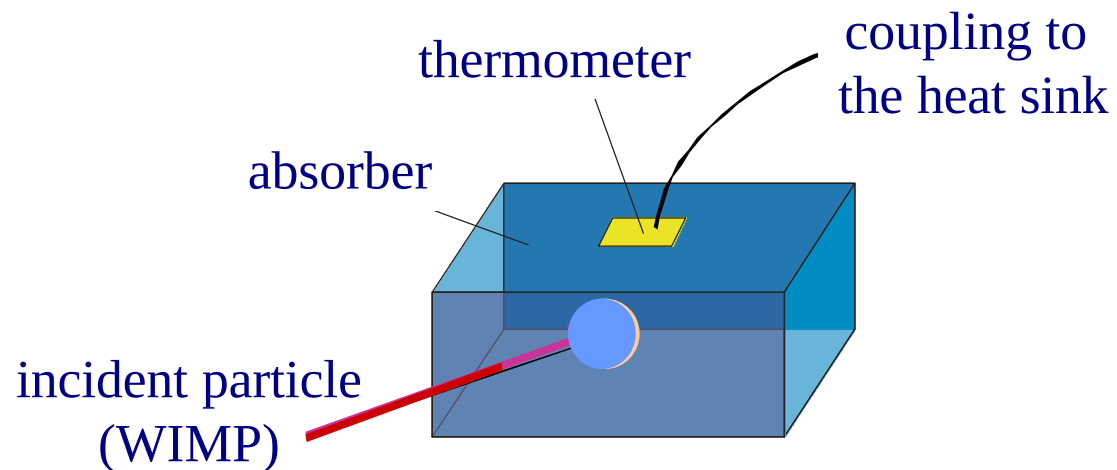


IPN Lyon
CEN Saclay
CNR Naples
CNRS Orsay
UC Berkeley
Yale University
Brown University
College de France
Oxford University
Stanford University
University of Tokio
Lancaster University
Universita di Milano
Paul Scherrer Institut
University of Twente
Universita di Genova
University of Wisconsin
European Space Agency
Moscow State University
Electronical Laboratory Japan
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Technische Universität München
Lawrence Livermore National Lab
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Eberhardt Karls Universität Tübingen
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Japan Atomic Energy Research Institute
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... ..

Cryogenic Calorimeters: Principles and Applications

Josef Jochum
Technical University Munich



Principle
Sensitivity - Resolution
Advantages - Applications

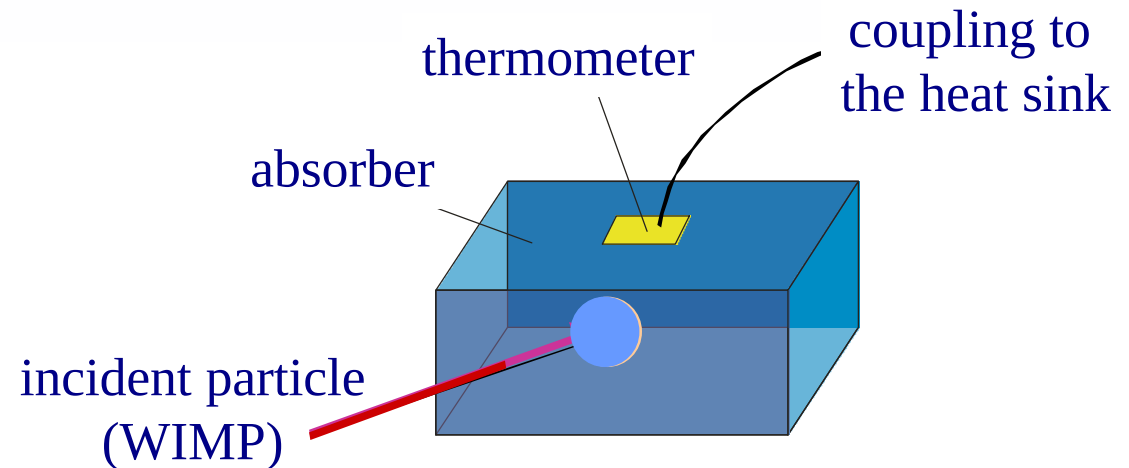
Basic Concept

absorption of energy due to incident particle or radiation causes a rise in temperature

$$\Delta T \propto E/C$$

very low temperature
=> high sensitivity, since C is small

types of thermometers: Superconducting phase transition thermometers (SPT)
NTD - Ge thermistors (highly doped semiconductors), ...



History of Cryogenic Particle Detection

1935 suggestion of low temperature calorimetry for particle detection

Simon

1949 alpha particles with phase transition thermometers

Andrews et al.

1974 danger of cosmic rays for low temperature experiments...

Niinikoski

1984 suggestion to measure low energy neutrinos by
coherent scattering with cryogenic detectors

Drukier, Stodolsky

1984 suggestion for double beta decay experiments

Fiorini, Niinikoski

1987 first conference on

Low Temperatur Detector for Neutrinos and Dark Matter, Munich
since then LTD's: Annecy, Gran Sasso, Oxford, Berkeley,
Beatenberg (Switzerland), Munich,
Dalfsen (Netherlands), ...Madison

Last proceedings in *Special Issue of Nucl.Instr. & Meth. A, Vol.444, # 1,2*

Motivations

solar neutrino detection (inverse beta decay)
target nuclei of interest can be incorporated

free choice of detector material

WIMP detection, coherent neutrino scattering

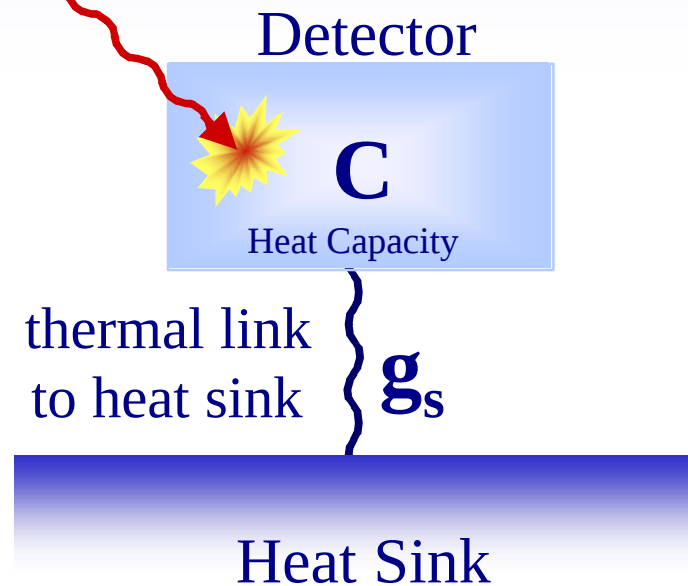
sensitivity to nuclear recoils

double beta decay, beta endpoint spectroscopy

source = detector
high energy resolution

Simplest Shape of the Signal

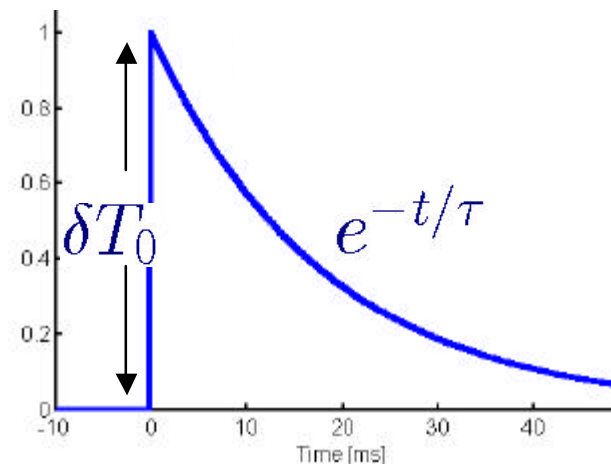
Radiation, Particle



initial temperature rise

$$\delta T_0 = \frac{E}{C}$$

relaxation to equilibrium
through thermal link



shape of the temperature signal

$$\delta T(t) = \frac{E}{C} e^{-t/\tau}$$

$$\tau = C/g_s$$

Heat Capacity

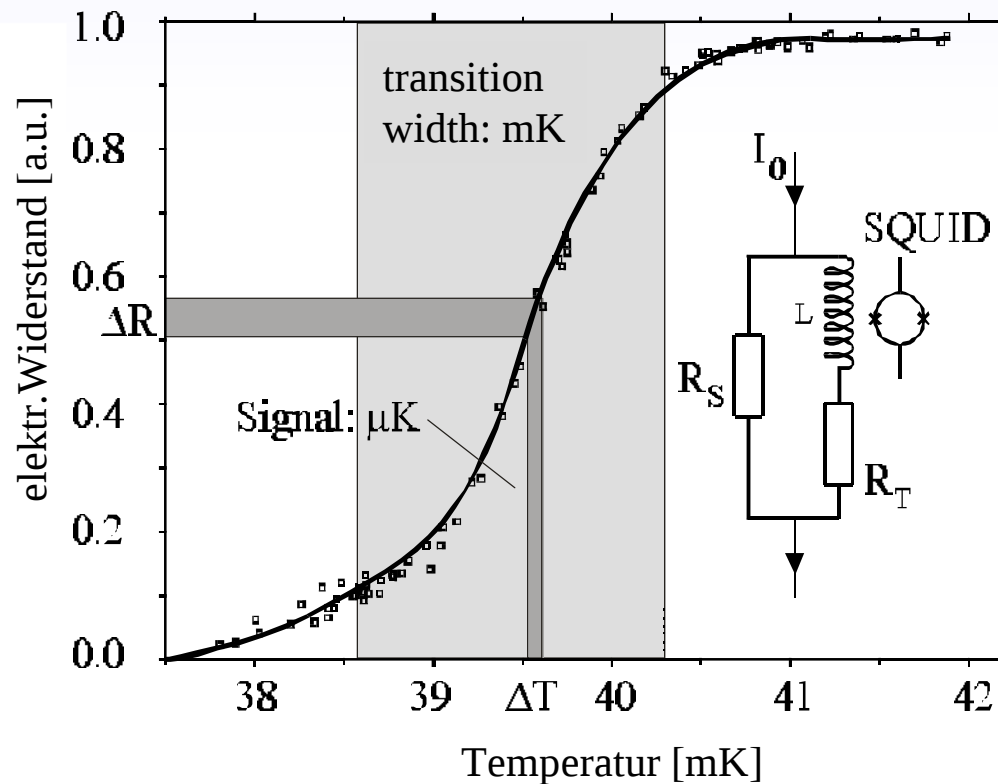
	T-dependence	material	heat capacity @25mK	temperature rise for 6keV @ 25mK
phonons	$C_{ph} \propto \left(\frac{T}{\Theta_D}\right)^3$	Saphir 1cm ³	5.5pJ/K = 34keV/mK	170μK
electrons	$C_e = \gamma T$	Gold 1cm ³	1.7mJ/K = 10TeV/mK	5.6 10⁻⁷ μK
		Gold (300μm x 300μm x 1μm)	170pJ/K = 1MeV/mK	5.6 μK

=> small detectors $\approx 0.1\text{mm}^2 \times 1\mu\text{m}$

=> low temperatures 10 - 100 mK

metals <=> dielectrics
small *large*

Phase Transition Thermometer



superconducting thin film
at the phase transition

materials:

W

$T_c = 15\text{mK}$

Ir/Au double layers

$T_c = (20-110)\text{mK}$

...

sensitivity

$$\alpha = \frac{T}{R} \frac{dR}{dT}$$

small change in temperature
=> large change in resistance
read out by a SQUID

Doped Semiconductor Thermistors

implanting dopands e.g. P or As into Si wafers

or

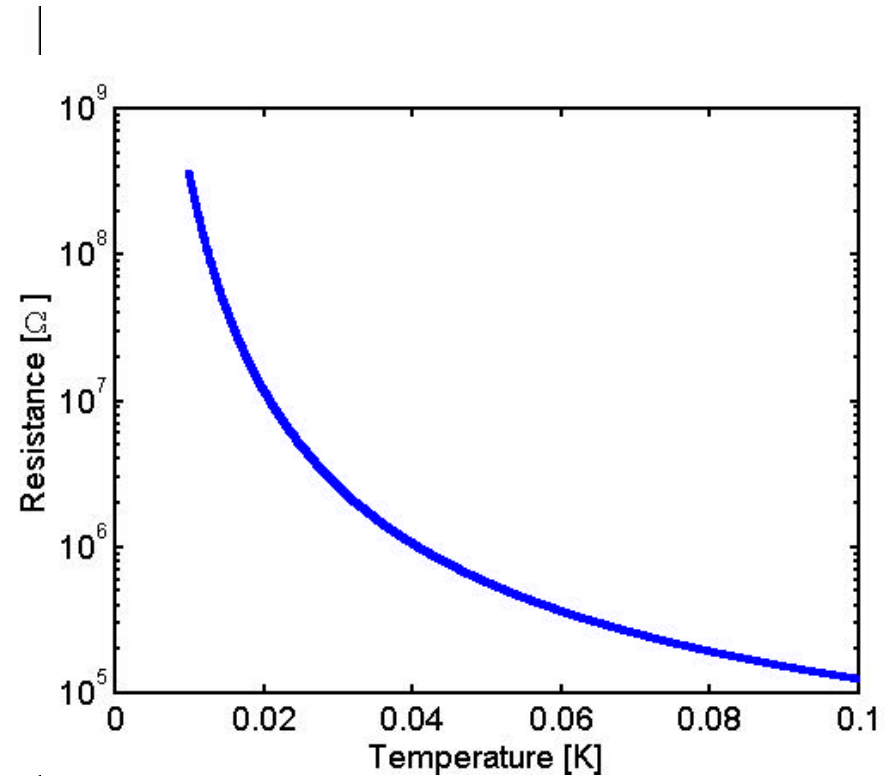
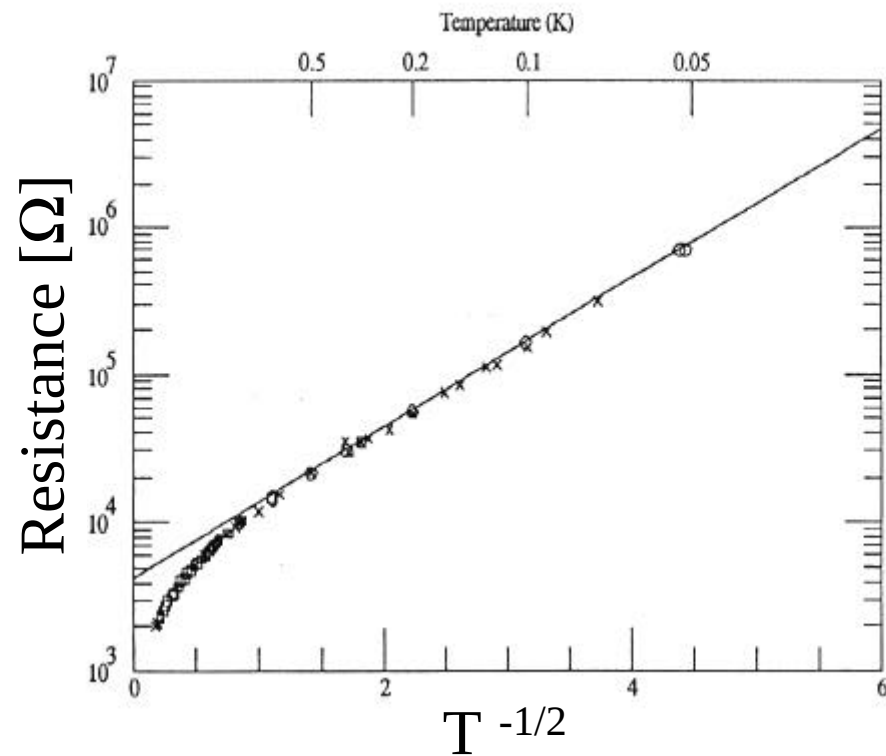
NTD Thermistors (uniform doping concentration by
irradiation of Ge to thermal neutrons,
Neutron **T**ransmutation **D**oping)

phonon assisted variable range hopping conduction
(carriers are localized at impurities and are aided
in tunneling by phonons)

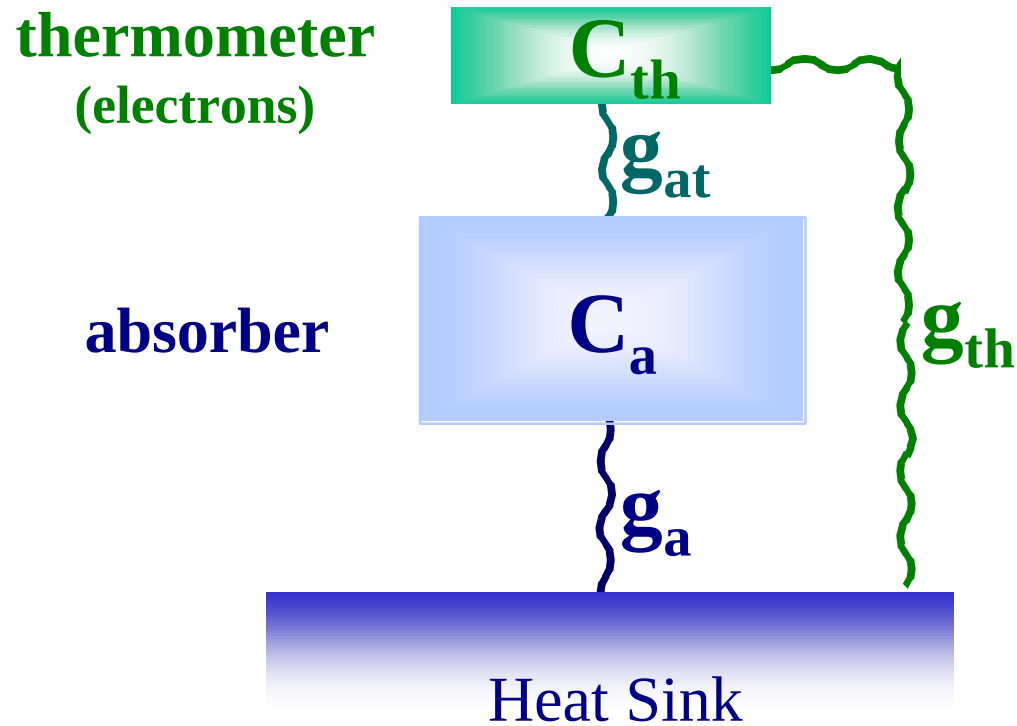
strong T dependence of resistance

$$R = R_o e^{\sqrt{T_o/T}}$$

Doped Semiconductor Thermistors



Calorimetry - Bolometry

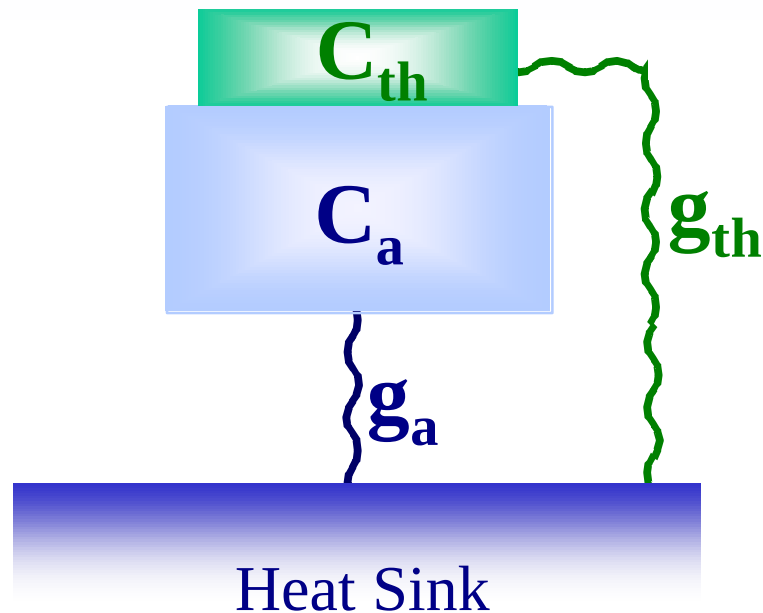


- g_{at} :
- contact resistance thermometer - absorber
 $\sim T^3$
 - electron-phonon coupling
 $\sim T^5$

resistance therm. - absorber

low for high frequency phonons
high for low frequency phonons

g_{at} large - T high (e.g. 100mK)

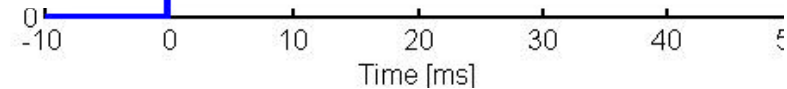


temperature
signal

$$E/C_{\text{tot}}$$

relaxation time

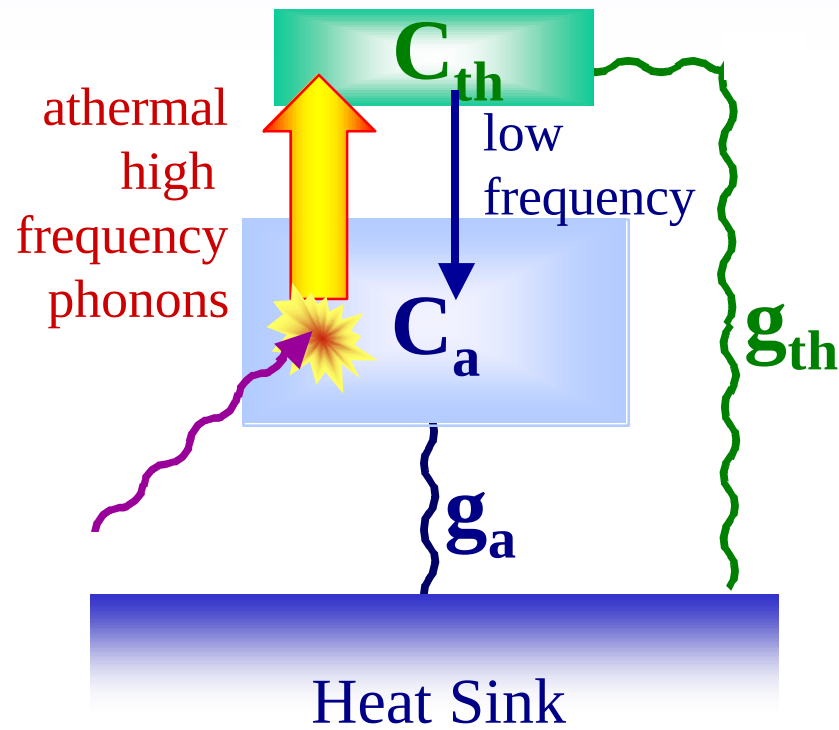
$$\tau = C_{\text{tot}}/g_{\text{tot}}$$



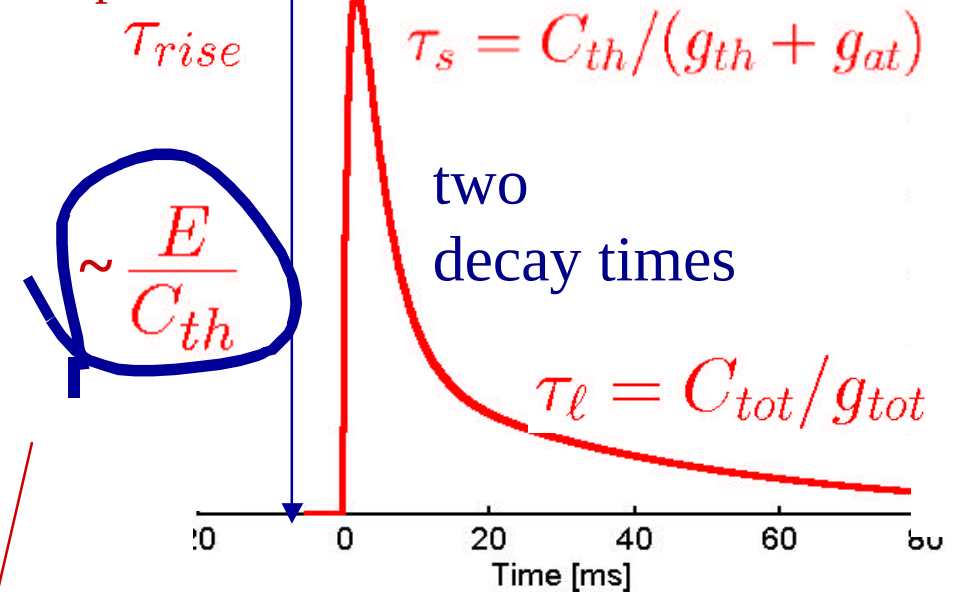
$$C_{\text{tot}} = C_a + C_{\text{th}} \quad g_{\text{tot}} = g_a + g_{\text{th}}$$

it's calorimetry

g_{at} small ($\begin{matrix} \text{high for low} \\ \text{low for high} \end{matrix}$ frequency) - T low (e.g. 15mK)

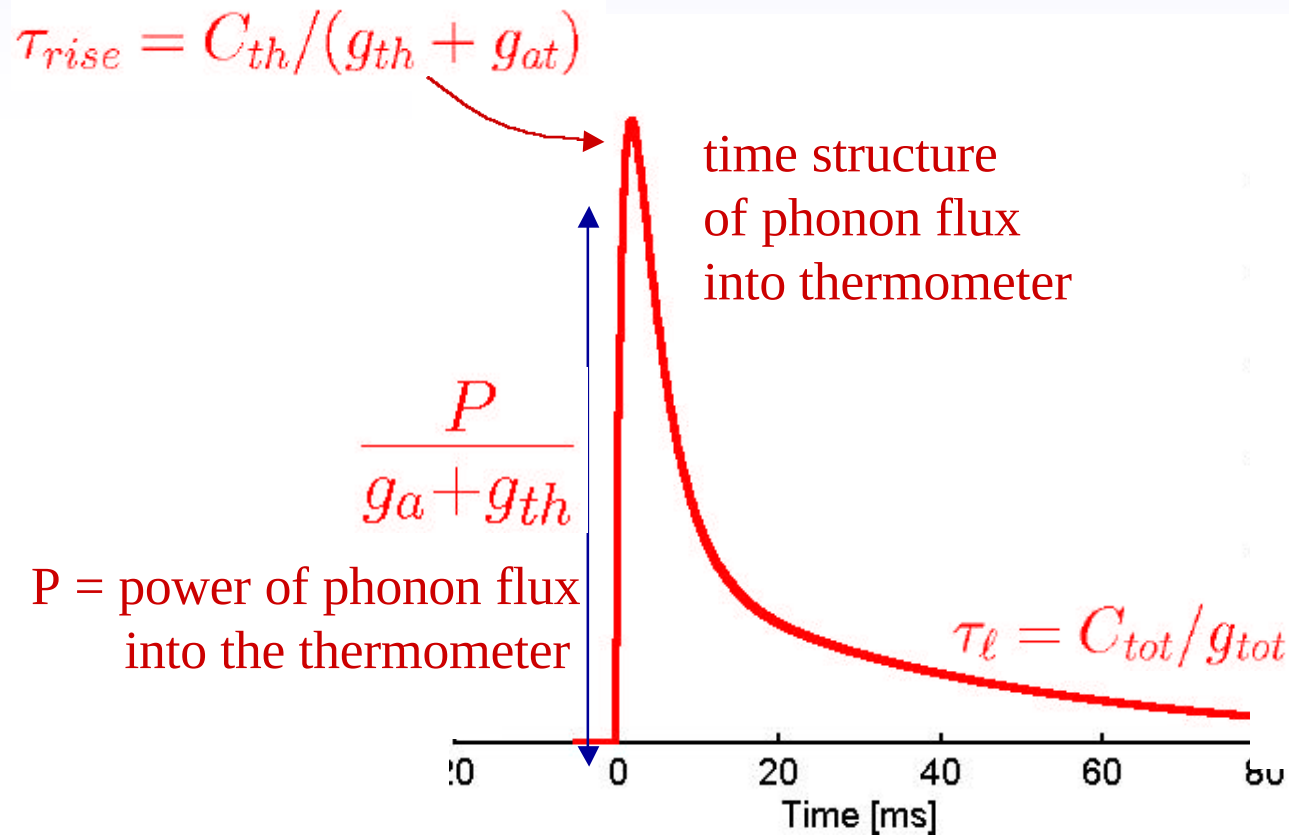


collection time
of high frequ.
phonons
 τ_{rise}



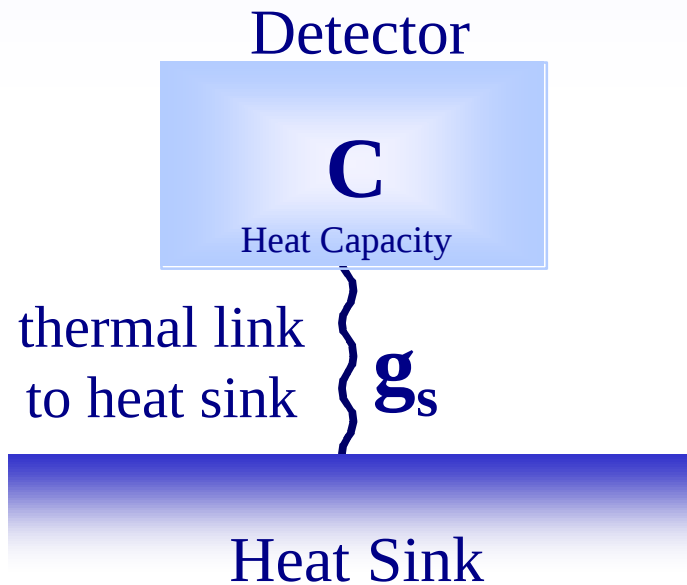
large absorbers possible !!!

Relaxation Time of Thermometer < Phonon Collection Time



fast signal **is bolometry** (for the phonons)
(later it's calorimetry again)

Sensitivity - Energy Resolution



statistical fluctuations
of energy content in C
due to random transport
of energy through g_s

=> noise spectrum

$$\frac{4k_B T g_s^2}{1 + \omega^2 \tau^2}$$

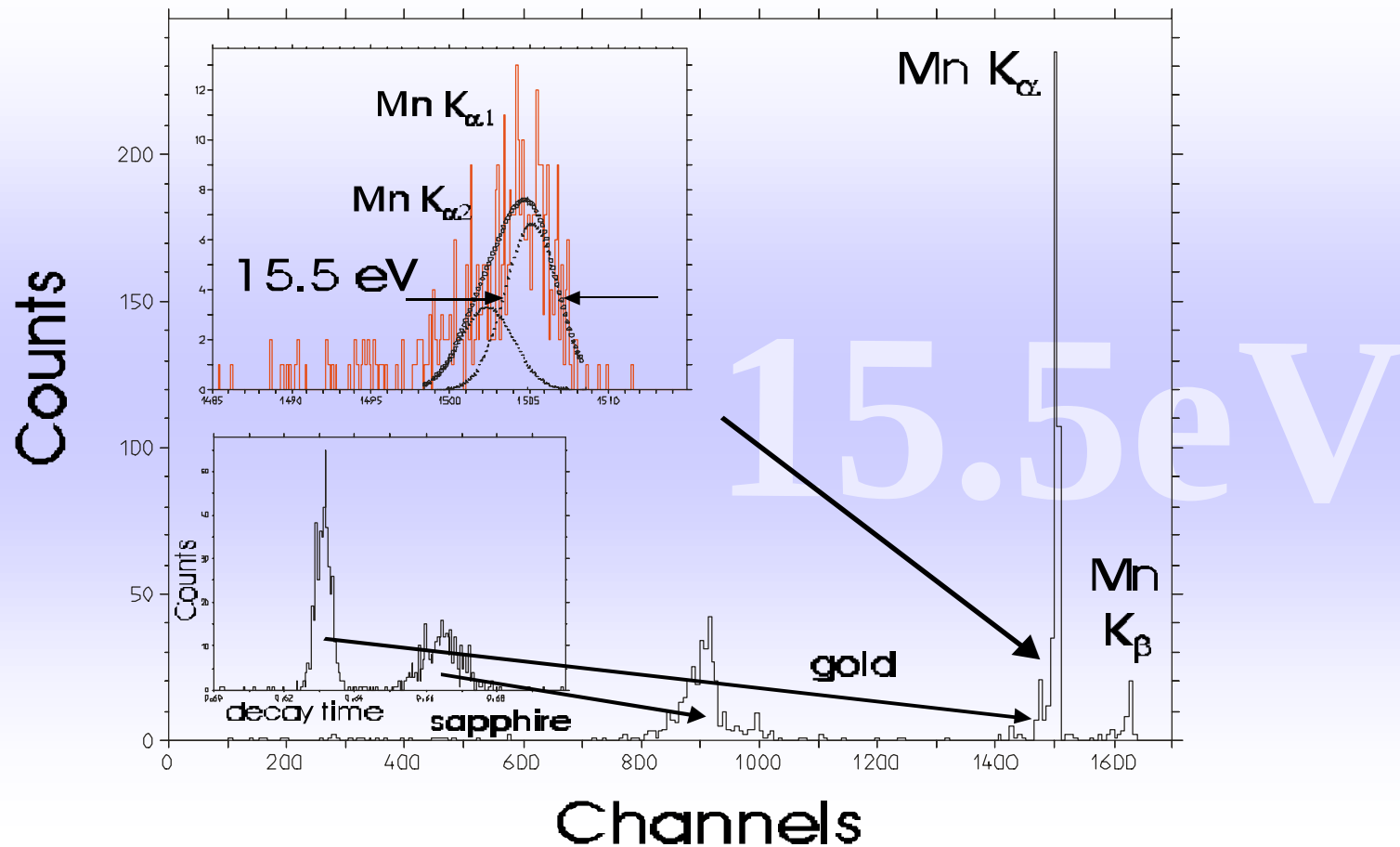
rms of energy content as 'limit' of energy resolution

$$\delta E^2 = 4kT^2C \quad \Rightarrow \quad \delta Signal \approx \sqrt{4kT^2C/\alpha}$$

measurement

Mikrokalorimeter:

Gold (300 μm x 300 μm x 1 μm) + Phase Transition Thermometer

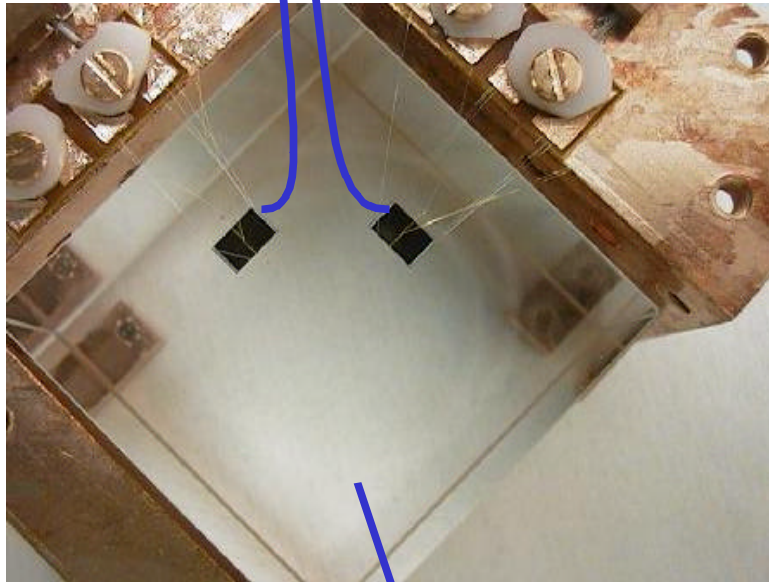


$$2.35\sqrt{4kTC^2} = 35\text{eV} \longrightarrow 1/\sqrt{\alpha} \text{ thermometer sensitivity}$$

CRESST Calorimeter

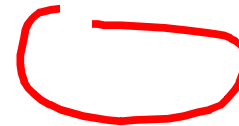
Tungsten Thermometer

Tc 15mK



Sapphire- Absorber
250gr, 4cm x 4cm x 4cm

energy threshold, a few 100eV !



$$2.35\sqrt{4kTC^2} = 52eV \rightarrow ???$$

Other Contributions to Energy Resolution -Band Gap

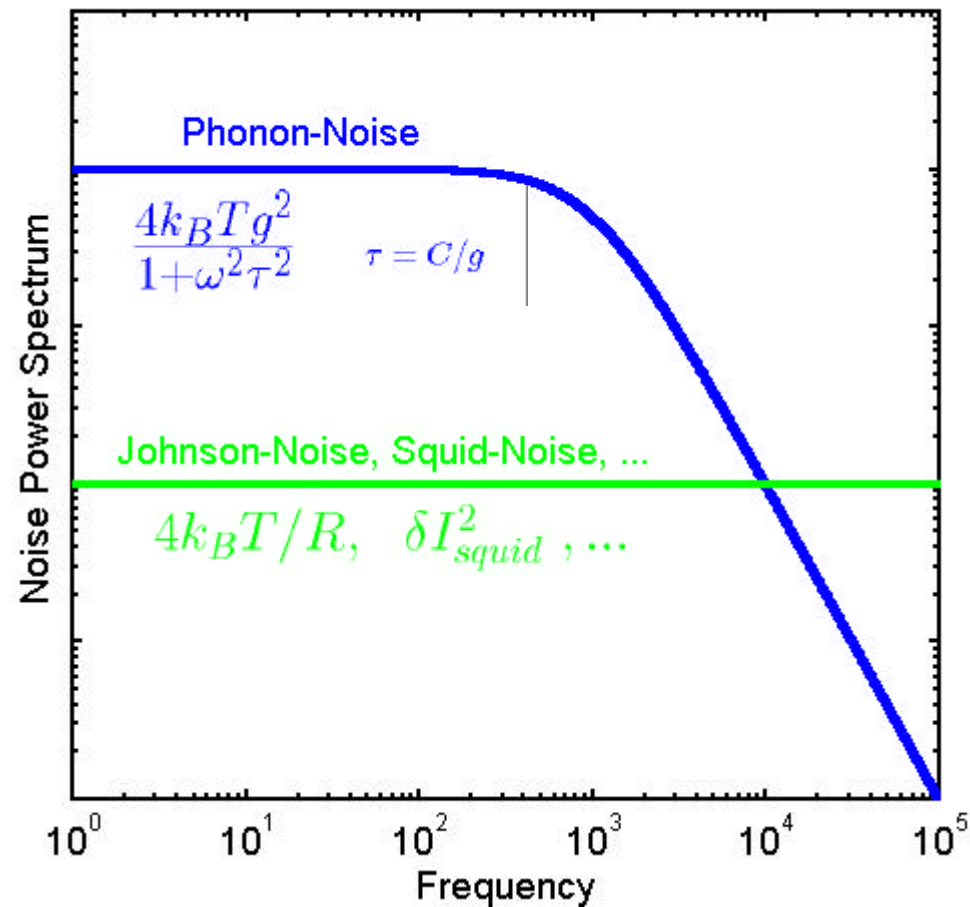
Statistical Limit: $\frac{\Delta E}{E} \propto \frac{1}{\sqrt{N}}$; $N = E / \epsilon$

Detector Material	effective excitation energy ϵ	number of excitations @ 1.5keV	energy resolution FWHM @ 6keV
Dielectric	10-20eV	75-150	$\approx 100 - 200$ eV
Semiconductor	3eV	500	≈ 30 eV
Superconductor	10^{-3} eV	$> 10^6$	< 1.5 eV

if excitations do not thermalize in a reasonable time => statistics

Other Contributions to Energy Resolution -Electronic Noise

- Johnson noise of thermometer resistance
- noise of preamplifier (SQUID, FET, ...)
- ...



Size - Sensitivity - Material

large detectors

high resolution

dielectrics
to get low C

metals
to avoid bandgap

reduced resolution

small detectors

Cryogenic Calorimeter: Advantages

very sensitive,
also to nuclear recoils

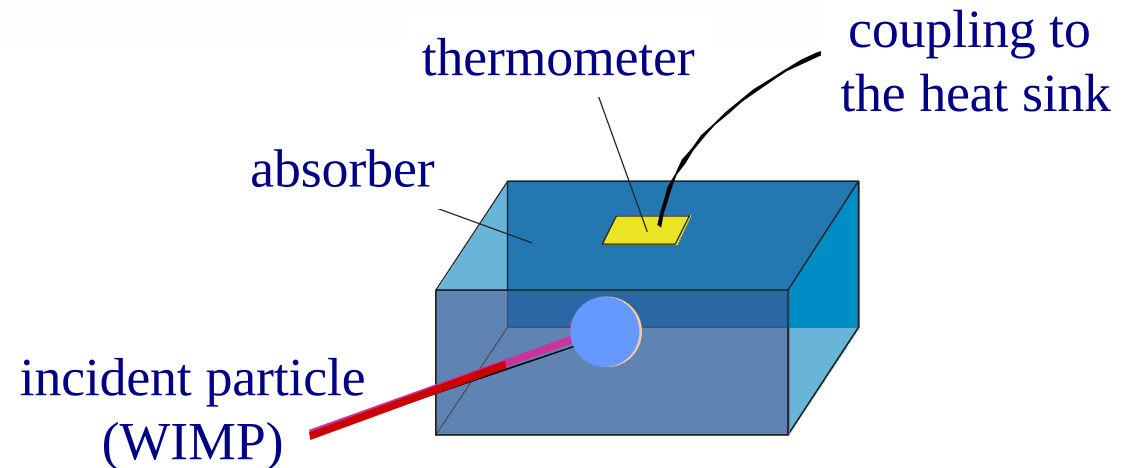
low energy threshold

high energy resolution

large variety of possible
target nuclei

very radiation hard

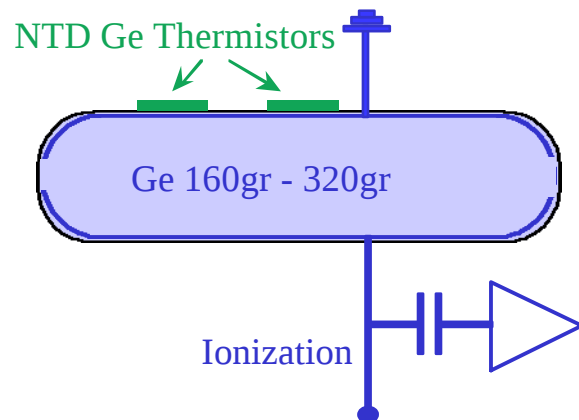
can be combined with other detection
methods (ionisation, szintillation)



$$\Delta T \propto E / C_{\text{Thermometer}}$$

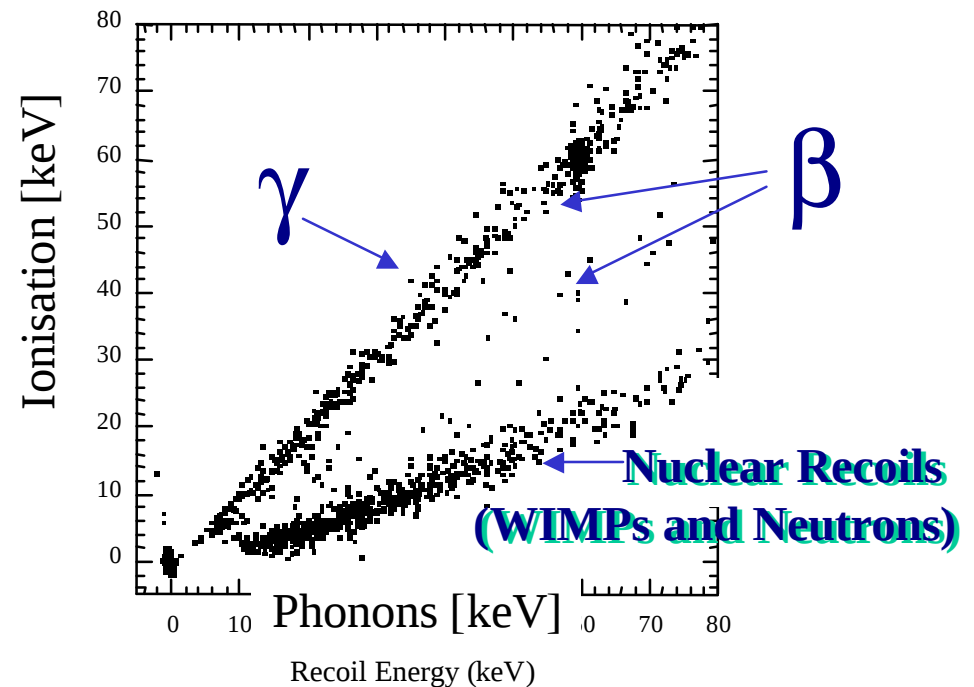
Discrimination: Photons $\leftrightarrow$ Nuclear Recoils

Simultaneous Measurement of Phonons and Ionisation



signal (nuclear recoils) and
dominant background (γ and β)
can be distinguished

CDMS and EDELWEISS Particle Dark Matter Searches



Low Temperature Detectors

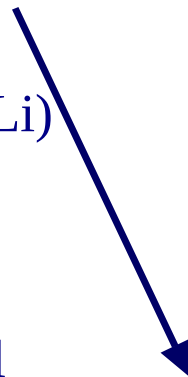
Mikrocalorimeter
Cryogenic-Calorimeter
Superconducting Tunneling Junctions
Magnetic Bolometers,
...

high energy resolution
(keV X-Ray: 10 x better than SiLi)

low energy threshold

free choice of detector material
(detector = source)

radiation hard



Microanalysis

Element-Analysis (low Z, PIXE)
Fluorescence-Analysis (XRF)
X-Ray Microscopes
X-Ray Emission in SEM

Optical- and X-Ray

Astrophysics

Next Generation after CCD
Energy-, Time- and Position-
Resolution for single Photons

Particle-Astrophysics

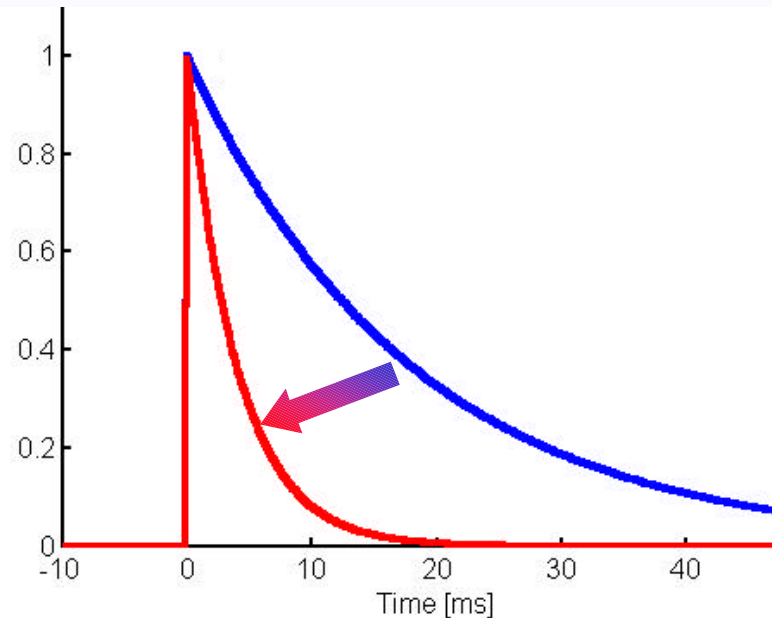
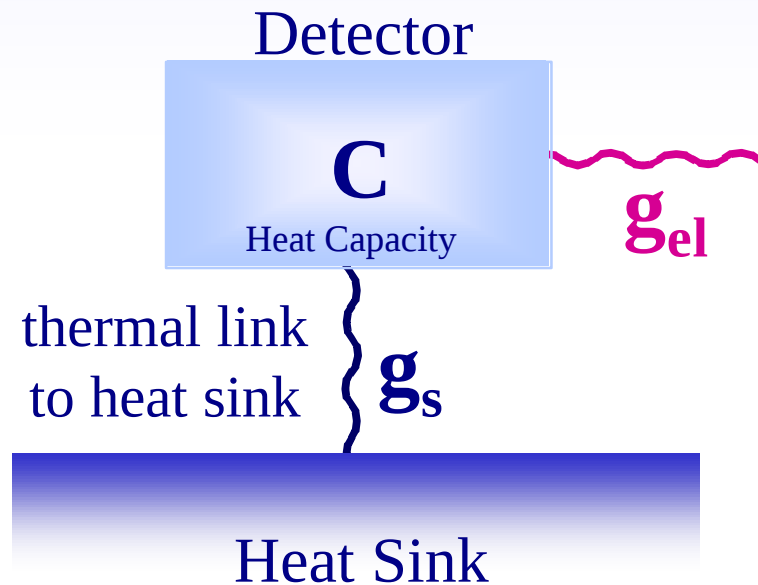
Nuclear Recoils / Dark Matter

Neutrino physics

β -decay und 2β -decay

Heavy Ion Physics

Heating Power by Measurement



if thermometer is voltage biased

$$P_{el} = U^2 / R_{thermometer}$$

if $T \nearrow \Rightarrow R_{thermometer} \nearrow$
(phase transition thermometer)

$$\Rightarrow P_{el} \searrow$$

speeds up relaxation to equilibrium

$$\text{relaxation time } \tau = \frac{C}{g_s + g_{el}}$$

$$g_{el} = -\frac{\delta P_{el}}{\delta T}$$